

## Durham Research Online

---

### Deposited in DRO:

27 February 2017

### Version of attached file:

Accepted Version

### Peer-review status of attached file:

Peer-reviewed

### Citation for published item:

Lin, Q. and Neethling, S.J. and Dobson, K.J. and Courtois, L. and Lee, P.D. (2015) 'Quantifying and minimising systematic and random errors in X-ray micro-tomography based volume measurements.', *Computers geosciences.*, 77 . pp. 1-7.

### Further information on publisher's website:

<https://doi.org/10.1016/j.cageo.2014.12.008>

### Publisher's copyright statement:

© 2015 The Authors. Published by Elsevier Ltd. This is an open access article under the CC BY license (<http://creativecommons.org/licenses/by/4.0/>)

### Additional information:

## Use policy

---

The full-text may be used and/or reproduced, and given to third parties in any format or medium, without prior permission or charge, for personal research or study, educational, or not-for-profit purposes provided that:

- a full bibliographic reference is made to the original source
- a [link](#) is made to the metadata record in DRO
- the full-text is not changed in any way

The full-text must not be sold in any format or medium without the formal permission of the copyright holders.

Please consult the [full DRO policy](#) for further details.

# Quantifying and minimising systematic and random errors in X-ray micro tomography based volume measurements

Q. Lin<sup>1</sup>, S.J. Neethling<sup>1\*</sup>, K.J.Dobson<sup>2,3</sup>, L.Courtois<sup>2,3</sup>, P.D. Lee<sup>2,3</sup>

<sup>1</sup>Rio Tinto Centre for Advanced Mineral Recovery, Department of Earth Science and Engineering, Imperial College London, London SW7 2AZ, United Kingdom

<sup>2</sup>Manchester X-ray Imaging Facility, School of Materials, University of Manchester, Oxford Rd., M13 9PL, UK

<sup>3</sup>Research Complex at Harwell, Rutherford Appleton Laboratories, Harwell, Didcot, Oxfordshire, OX11 0FA, UK

\*Corresponding author:

Tel.: +44 2075949341.

E-mail address: [s.neethling@imperial.ac.uk](mailto:s.neethling@imperial.ac.uk) (S.J. Neethling)

Address: Department of Earth Science and Engineering, Royal School of Mines, Imperial College London

Postcode: SW7 2BP

Keywords: Micro-CT, grain tracking, uncertainty, error propagation

## Abstract

X-ray micro tomography is increasingly used for the quantitative analysis of the volumes of features within the 3D images. As with any measurement, there will be error and uncertainty associated with these measurements. In this paper a method for quantifying both the systematic and random components of this error in the measured volume is presented. The systematic error is the offset between the actual and measured volume which is consistent between different measurements and can therefore be eliminated by appropriate calibration. In XMT measurements this is often caused by an inappropriate threshold value. The random error is not associated with any systematic offset in the measured volume and could be caused, for instance, by variations in the location of the specific object relative to the voxel grid. It can be eliminated by repeated measurements. It was found that both the systematic and random components of the error are a strong function of the size of the object measured relative to the voxel size. The relative error in the volume was found to follow approximately a power law relationship with the volume of the object, but with an exponent that implied, unexpectedly, that the relative error was proportional to the radius of the object for small objects, though the exponent did imply that the relative error was approximately proportional to the surface area of the object for larger objects. In an example application involving the size of mineral grains in an ore sample, the uncertainty associated with the random error in the volume is larger than the object itself for objects smaller than about 8 voxels and is greater than 10% for any object smaller than about 260 voxels. A methodology is presented for reducing the random error by combining the results from either multiple scans of the same object or scans of multiple similar objects, with an uncertainty of less than 5% requiring 12 objects of 100 voxels or 600 objects of 4 voxels. As the systematic error in a measurement

cannot be eliminated by combining the results from multiple measurements, this paper introduces a procedure for using volume standards to reduce the systematic error, especially for smaller objects where the relative error is larger.

## 1. Introduction

X-ray micro tomography (XMT) is a popular technique for the non-destructive qualitative and quantitative investigation of the internal structure of objects. It has been widely applied across material science (Puncreobutr *et al.*, 2012; Stock, 1999), engineering (Aydoğan *et al.*, 2006); (Ghorbani *et al.*, 2011; Ketcham and Carlson, 2001) and biological sciences (Yue *et al.*, 2011) to provide quantitative data about the structure and morphology of 3D objects and features within them (crystals, pores, fractures etc.).

For the measurement of object or feature volumes from XMT images, each voxel (smallest volume element, equivalent to a 3D pixel) belonging to a feature or object is obtained using a thresholding algorithm, and the volume obtained by counting the relevant voxels. However, the boundaries of features rarely coincide with the boundaries of the regular voxel grid, leading to the “partial volume effect” at the interface, where voxels have an intermediate bulk composition, and there is some uncertainty in the exact boundary location (Ketcham and Carlson, 2001; Stock, 1999). Theoretically the partial volume effect should only affect a narrow (few voxel) region when the boundary is planar, smooth and sharp, but in certain systems boundaries can be uneven and diffuse (Figure 1).

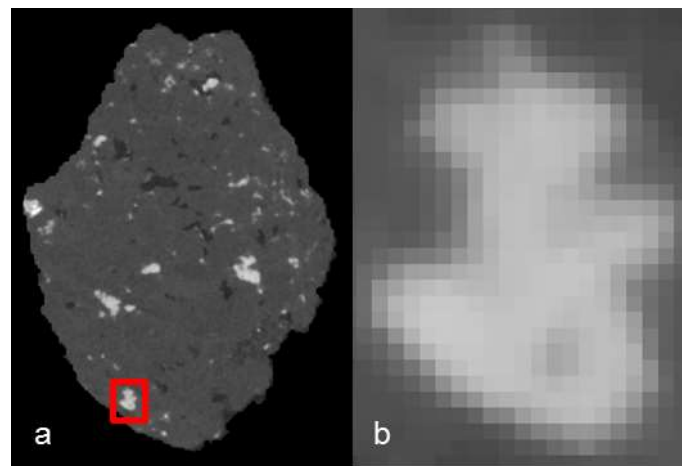


Figure 1. a) A 2D slice through a 3D tomography volume of example data showing mineral grains within an ore particle. b) A region of interest demonstrating diffuse boundaries.

The choice of thresholding algorithm or threshold value will have a systematic effect on the measured volume, while variability in the exact location of the object relative to the voxel grid will cause a random variation in the measured volume. The relative impact on the measured volume of both these systematic and random errors will be strongly dependent on the size of the object relative to the voxel size, as the proportion of the volume that is within the uncertain region at the boundaries of objects will decrease as the object size increases.

64 In this paper we describe a procedure for quantifying both the systematic and random  
65 components of this uncertainty in volume. In particular, we describe how to ascertain how many  
66 times an object needs to be scanned (or how many similar objects in the same scan need to be  
67 combined) to achieve a given level of accuracy in the measured volume, assuming that any  
68 systematic error has been eliminated. Repeatability will also be influenced by both the random  
69 and systematic components of the error as the systematic error is likely to change from scan to  
70 scan, while the random component will add uncertainty to the measurement. Although the  
71 methodology presented significantly improved repeatability, for absolute dimensional accuracy  
72 calibration with an appropriate phantom is required.

73 While our methodology is applicable to a wide range of 3D image analysis applications, the results  
74 obtained will depend to some degree on the sample being studied and the specifics of the scanner  
75 used. In this paper the example used is the quantification of metal sulphide grain volumes within  
76 an ore particle/rock fragment. The ore particles were scanned using a Nikon Metris Custom Bay  
77 with a 1 mm aluminium filter to reduce the effect of beam hardening, 89 kV energy, 0.708 s  
78 exposure time and 2001 projections. The detector size was 2000×2000 pixels, giving a linear  
79 resolution of approximately 17 microns for the magnification selected. We chose this example as  
80 there are a large number of mineral grains within the image volume and these grains are known to  
81 have a wide volume distribution. For the scan resolution used the mineral grains range from sub-  
82 voxel sizes to tens of thousands of voxels, allowing for the effect of the volume of the object to be  
83 studied over many orders of magnitude.

84 A key requirement of this methodology is the ability to identify the same objects in repeated scans.  
85 An algorithm developed for tracking the dissolution of mineral grains as they undergo leaching is  
86 used for this purpose. The first section of this paper thus gives a short description of this algorithm  
87 as the data generated from it is the source of the statistical analysis.

## 88 **2. Grain tracking and identification methodology**

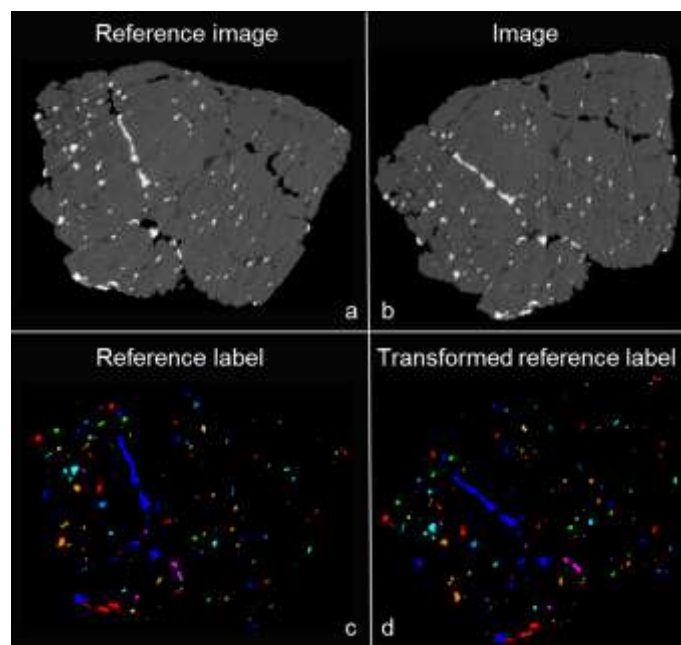
89 The procedure for the image processing was:

- 90 1. A 3×3×3 median filter was applied to reduce the noise level.
- 91 2. The transformation matrix to align subsequence scans to the orientation and location of  
92 the reference scan was calculated and extracted (Studholme *et al.*, 1999).
- 93 3. The threshold for distinguishing the ore particles from the air phase was obtained using the  
94 Otsu algorithm (Otsu, 1979), while the metal sulphide grains are distinguished from the ore  
95 matrix using a maximum entropy algorithm (Kapur *et al.*, 1985). The reason for the  
96 different algorithms is that the air and rock have very distinct peaks in the intensity  
97 histogram, while the relatively small volume of metal sulphide present means that there is  
98 no distinct peak in the histogram.
- 99 4. The individual grains were then tracked across different images.

100 The algorithm starts by identifying all the mineral grains of interest in the reference image. The  
101 connectivity of the grains are analysed so that each isolated grain is given a unique identifier. On  
102 subsequent images voxels that are identified as mineral grains need to be given the same  
103 identifier number as they had in the original image. This is achieved by using a mask based on the  
104 reference image. This mask is rotated and translated to match the location and orientation of the  
105 ore particle in the subsequent image. This mask is then applied to the mineral grains.

106 Since the grains do not grow between images, this masking would be all that is required if the  
107 thresholding of the images and the translation and rotation of the mask were perfect. In general  
108 this is not the case and unassigned rims can remain around the masked grains. This problem is  
109 resolved by assigning these rim voxels the identifier of a neighbouring identified voxel. This  
110 process is repeated until all voxels in the intensity range are identified or discarded.

111 It should be noted that in this algorithm it is the mask that is rotated and translated and not the  
112 data itself. Rotating the data would have an effect on the measured volume of the grains and thus  
113 also the error associated with the volume measurement as the interpolation required to project  
114 the rotated and translated data back onto a grid will cause the boundaries to become even more  
115 diffuse. Translating and rotating the mask will cause slight changes in the size and shape of  
116 masked regions, but this will have virtually no impact on the algorithm as the rims that result from  
117 slight errors in the mask are accounted for in the algorithm. Figure 2 shows an example of a  
118 reference and subsequent image as well as the original and transformed mask. Note that, while  
119 the figure shows a 2D slice, the rotations and translations were all 3D.



120  
121 **Figure 2** Example of reference label transformation. a) Reference image. b) Image in subsequence scan. c) Label mask for  
122 reference image. d) Transformed reference label after applying 4×4 transformation matrix. The transformation matrix is  
123 calculated using (a) and (b).

124 This identification method has a few assumptions and limitations. Firstly, any objects that do not  
125 appear in the initial image but exist in a later scan are not counted. This issue can occur for objects  
126 that are of a size very close to the voxel resolution or due to phantom particles caused by noise in  
127 the image, which can be ameliorated by the use of a median filter. Another potential issue with  
128 this algorithm is if the mask does not overlap any portion of the object in subsequent images.  
129 Again this is only likely for objects that are approximately the same size as the voxel resolution.  
130 Objects that appear in the reference scan, but are not observed in the subsequent scan are  
131 included in the statistics, though objects that are not in the reference scan, but appear in a  
132 subsequent scan are not counted. These objects make up about 5% of the total number of objects  
133 in the subsequent scan, but as their sizes are all close to the scan resolution they account for only  
134 0.05% of the total volume of the identified objects.

### 135 **3. Error and uncertainty in the volume of scanned objects**

136 Before the volume data can be used with confidence the systematic and random errors in the  
137 measurement need to be understood. Systematic errors are those in which the error is the same  
138 for all similar objects and, for volume measurement, will typically be a function of the size of that  
139 object. Correction of systematic errors is possible using appropriate standards and calibration.  
140 Random errors are those that are not the same for similar objects or between scans and thus add  
141 an uncertainty to measurements that cannot be eliminated by calibration. However, unlike  
142 systematic errors, random errors do not influence the average measured volume if enough  
143 volume measurements have been used. What this paper will demonstrate is a methodology for  
144 determining how many repeat measurements (or measurements of similar objects) need to be  
145 made to reduce the uncertainty caused by the random error to an appropriately small value (what  
146 is considered appropriately small will, of course, depend upon the application).

147 The systematic error in the grain volume will come about from effects such as an error in the  
148 threshold used, while the random error will come about due to effects such as the change in the  
149 partial volume effect due to the specific location of the mineral grain relative to the voxel grid,  
150 which will change from scan to scan and from grain to grain.

#### 151 **3.1 Sensitivity of measured volume to threshold changes**

152 Global thresholding is a common method to distinguish different phases (Gonzalez *et al.*, 2003),  
153 and the choice of threshold used to distinguish the phases can have a large effect on the volume  
154 measured. Thresholding is an important step in the image quantification and it has a direct  
155 relationship with the uncertainty, especially for smaller grains. Much of this uncertainty arises  
156 from partial volume effect, where the edges of grains are blurred due to the fact that they do not  
157 necessarily align with the voxels. Typically an algorithm is used to choose the threshold to reduce  
158 subjectivity in the identification of the objects within the image, but this does not mean that  
159 systematic errors due to thresholding are eliminated, though, and it might well be appropriate to

160 adjust the threshold value to minimise these systematic errors if an accurate, rather than simply  
161 consistent, volume measurement is required.

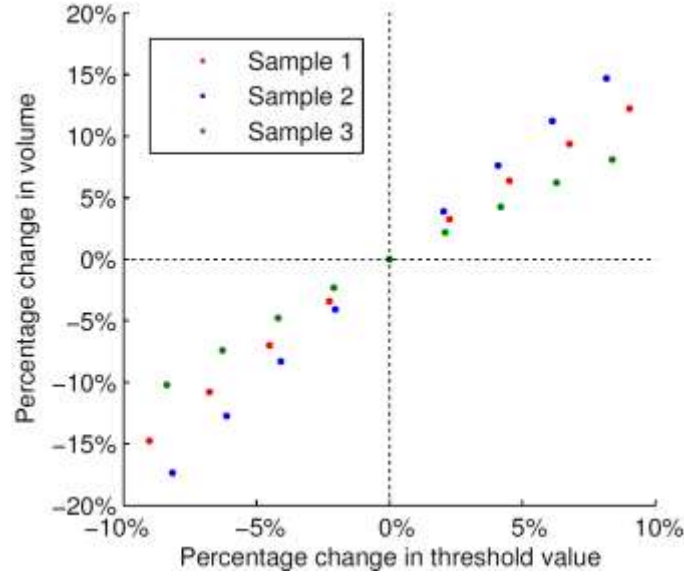
162 Local thresholding algorithms can also be used (Gonzalez *et al.*, 2003). While the trends in random  
163 error associated with local algorithms are likely to be similar to those associated with global  
164 methods, as these errors are largely associated with the real uncertainty in the images, the  
165 systematic errors will be very algorithm specific. For this reason this paper concentrates on the  
166 uncertainties and errors associated with global thresholding as these responses are the same  
167 irrespective of which algorithm is used to choose the threshold (the response in the measured  
168 volume brought about by varying the threshold value will be the same irrespective of the  
169 algorithm used to obtain the initial threshold value).

170 The initial threshold values used to identify the mineral grains were obtained by applying the  
171 maximum entropy global thresholding algorithm to each rock (Kapur *et al.*, 1985). The threshold  
172 was then adjusted from these values and the percentage change in the measured total volume of  
173 all the mineral grains calculated (Figure 3). The shift in threshold value is quantified using the ratio  
174 between absolute shift in the value and the difference between the rock and mineral grain phase  
175 thresholds:

176 
$$T_{shift} \% = \frac{T_{shift}}{T_{grain} - T_{rock}} \quad (1)$$

177 where  $T_{grain}$  is the threshold for the sulphide grains, and  $T_{rock}$  is the threshold for rock phase. The  
178 reason for using the change relative to this difference is that the appropriate threshold value must  
179 lie between the intensity of the grains and the matrix.

180 There is an approximately linear variation in the measured volume as the threshold value is  
181 changed (Figure 3), though the magnitude of the variation changes somewhat from sample to  
182 sample. This variability is probably due to differences in the size distribution of the grains within  
183 the three rocks.



**Figure 3. The relationship between the change in mineral grain volume and the variation in the threshold value**

The sensitivity of a grain's measured volume to a change in threshold is very dependent on their size relative to the voxel resolution. Smaller grains are more sensitive to a change in threshold because this is mainly a surface effect and smaller grains have a larger specific surface area. Assuming that a small change in the threshold produces a small change in the location of the boundary (this analysis does not require that the relationship between the change in the position of boundary and the threshold be a simple one, only that the change in position is approximately the same at all boundaries), the fractional change in volume can be expressed as:

$$\frac{\Delta V}{V} \propto \frac{r^2 \Delta r}{r^3} = k V^{-\frac{1}{3}} \Delta r \quad (2)$$

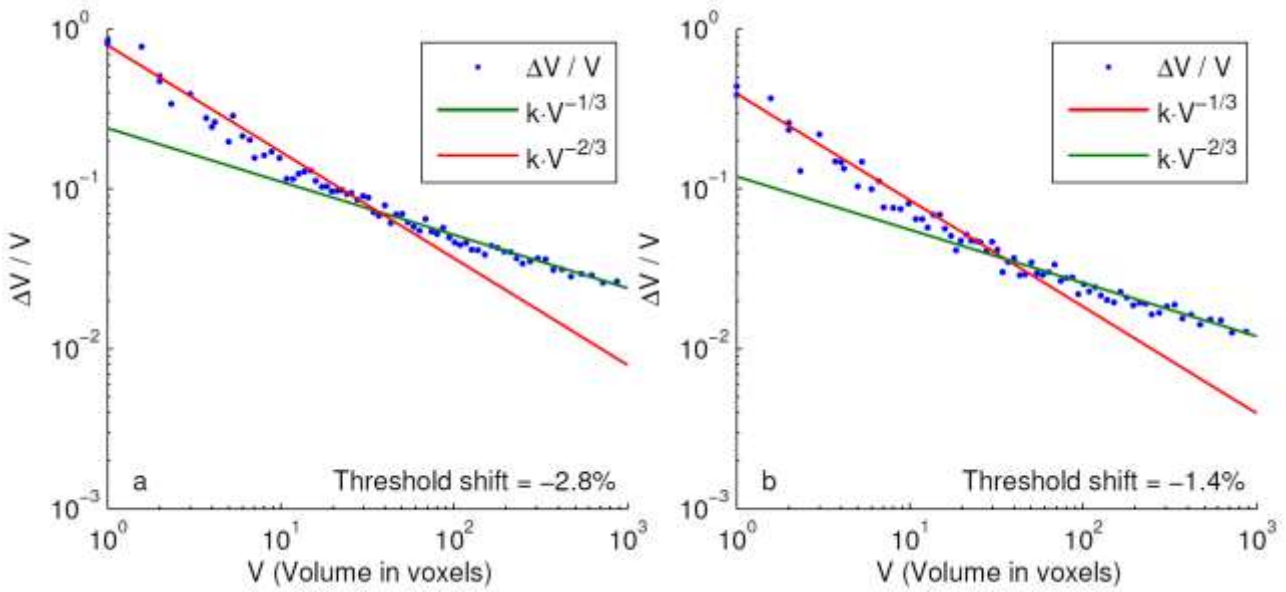
where  $V$  is the volume of the grain,  $r$  is a linear dimension of the object (proportional to  $V^{1/3}$ ) and  $k$  is a dimensionless constant.  $\Delta r$  is the change in the position of the boundary, which mainly depends on the change in the threshold value, but can also depend on the shape and size of grains. A power law exponent of  $-1/3$  implies that the relative change in volume is inversely proportional to the grain radius and proportional to its specific surface area.

Plotting  $\Delta V/V$  against grain volume for different threshold (Figure 4) values shows that the larger grains ( $> \sim 35$  voxels) follow Equation (2), but that the smaller grains ( $< \sim 35$  voxels) have a more negative slope, with  $\Delta V/V = k V^{-2/3}$  (a power law exponent of  $-2/3$ ) producing a better fit. An exponent of  $-2/3$  implies that the change in the volume upon a threshold change scales with the radius of the grains rather than its area for the smaller grains. This is somewhat unexpected, though one possible explanation for this is that either the reconstruction algorithm or the imaging itself is producing more uncertainty in one of the axes than the others. Another possible explanation is that the apparent shape of some of the objects are strong functions of the threshold value chosen. This is not much of an issue for convex objects, but is likely to be



208 important for more complex objects. For such objects simple thresholding might not be sufficient  
 209 and more complex techniques may need to be applied. An example of such a method is de-  
 210 convolution based on the assumption that the blurring of the edges takes the form of a point  
 211 spread function (Ketcham, 2006; Ketcham and Hildebrandt, 2014).

212 While the relative change in the average measured grain volume is of a similar magnitude to  $T_{shift}$   
 213 (Figure 3), for individual grains the difference is strongly size dependant (Figure 4). Since the small  
 214 grains are more sensitive to changes in threshold than the larger grains, it is this region of the  
 215 curve that is most important.



216  
 217 **Figure 4.** The plot of the relative change in grain volume as a function of the volume for two different threshold changes. a) 2.8%  
 218 b) 1.4%. The power law relationships for large ( $-1/3$ ) and small grains ( $-2/3$ ) are also shown.

219 In Figure 5, the prefactor in the best fit to smaller grains (less than 100 voxels) for a power law  
 220 relationship with an exponent of  $-2/3$  is plotted against the change in the threshold value. Since  
 221 the  $k$  value and the magnitude of the average change in volume are directly related, there is also a  
 222 near linear change in  $k$  with the change in the threshold value. This curve will be used later to  
 223 correct the systematic errors (Section 4).

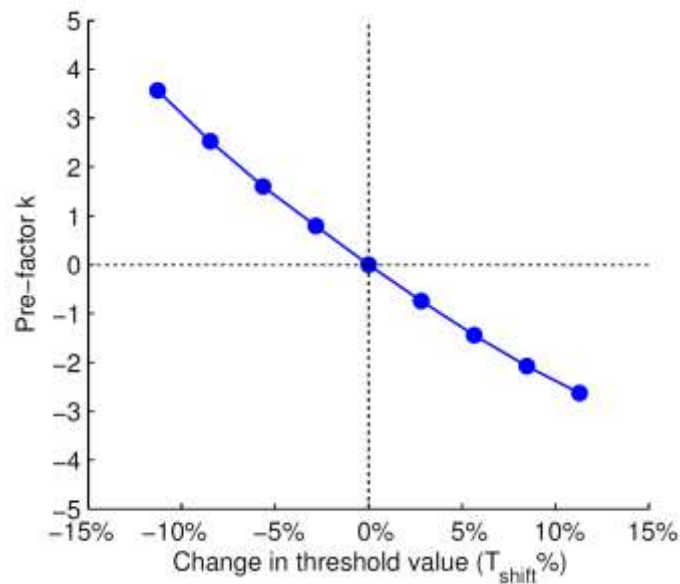


Figure 5. Plot of prefactor k as threshold value changes

### 3.2 Estimation of grain volume uncertainty

While the effect of changing the threshold can be obtained from a single image, repeat scans of the same volume are required to determine the random component of the error. As the scanned volume contains a number of ore particles and each ore particle contains thousands of grains, the identification procedure outlined in Section 2 allows us to look at the variability in the measured volume of tens of thousands of individual objects. The same analysis can be carried out for systems containing fewer objects, but in order to generate sufficient statistics on which to base the analysis, repeated scans of the same objects may need to be carried out.

Taking two images of the same sample volume, the relative error in volume measurement for each individual grain was calculated. The grain volumes were then ordered according to size and the standard deviation in the relative error was calculated for sets of 500 grains of similar volume and plotted against the mean volume of the set of grains (Figure 6).

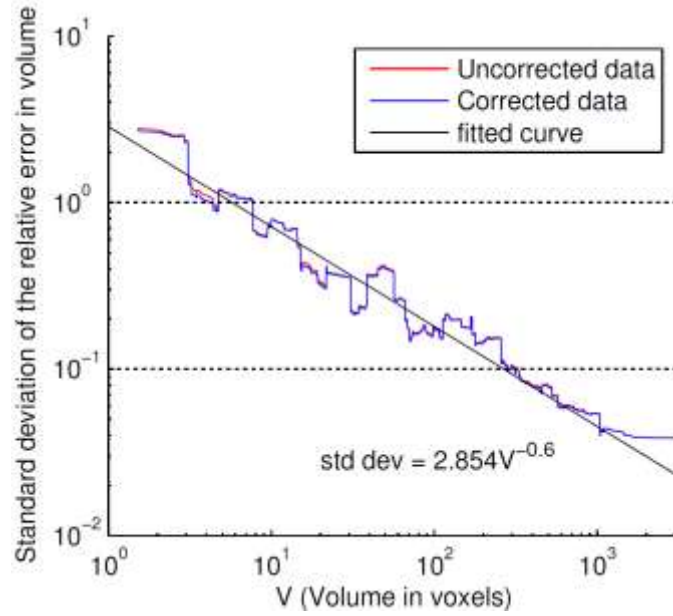


Figure 6. Standard deviation in the relative error in the grain volume as a function of the grain volume

For a single grain the uncertainty (expressed as the standard deviation of the measurement) in the size of the grain is as large as the grain itself for any grain less than approximately 8 voxels. For an uncertainty of less than 10% the grain needs to be larger than about 260 voxels in volume. Figure 6 shows that there is a power law relationship between the standard deviation in the relative error and grain volume, with an exponent of close to  $-2/3$ , which is consistent with the scaling for the systematic error<sup>1</sup> (Figure 4). This means that the magnitude of the random component of the error is approximately proportional to the radius of the grain, which is again surprising as the naive expectation would be that this error would be related to the surface area of the grain.

The uncertainty in the measured volume can be reduced by either repeated scans of the same object or by combining the results from a number of similar objects. As the uncertainty for an individual object is a function of the object volume, the number of similar objects,  $N$ , of volume  $V$  that need to be combined to achieve an acceptable relative error,  $\varepsilon$ , in the measured volume can be calculated (or, alternatively,  $N$  is the number of times that the same object needs to be scanned):

$$N = \frac{(\kappa V^n)^2}{\varepsilon^2} \quad (3)$$

where  $\kappa$  is the prefactor in the relationship between the relative standard deviation in the measure volume of a single grain and  $n$  is the power law exponent. For example, based on the

---

<sup>1</sup> While they have similar volume scalings in this system, there is no fundamental reason why the systematic and random components of an error need to have the same dependencies.

scans used to produce Figure 6, to reduce the random component of the uncertainty when measuring volume to less than 5% you would need to combine the measurements from approximately 600 grains with a volume of 4 voxels, approximately 12 grains with a volume of 100 voxels objects, or one object of 1000 voxels (see Figure 7). This will not account for any genuine variability in the behaviour of nominally identical objects, and it is important to note that it is only the random component of the error that is reduced by averaging repeat results. By definition, combining results will have no impact on any systematic error. Figure 7 can only be used as an indication of the error expected as the error will depend upon the particular material and its scanning conditions.

The procedure presented is relatively straight forward, and is recommended whenever precise quantitative data for the volume or volume change is required.

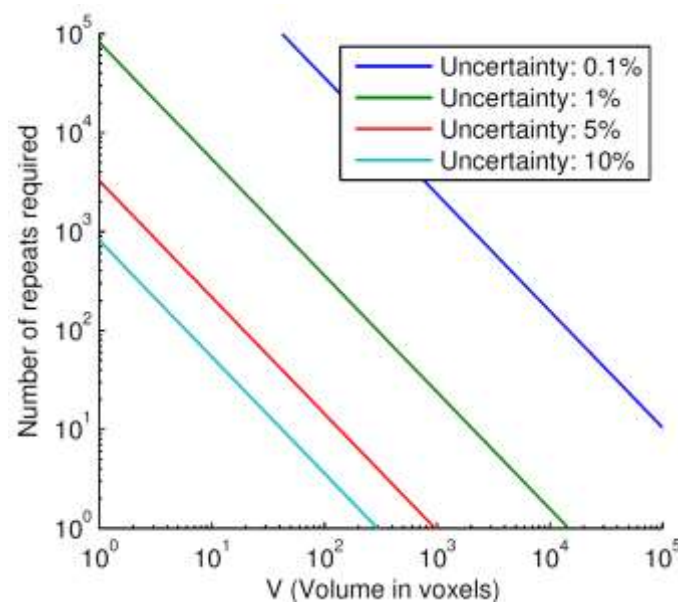


Figure 7. Number of repeats required to reduce the random component of the uncertainty (relative standard deviation) in the volume measurement to a given level as function of the object volume.

#### 4. Obtaining consistent results in the face of systematic errors

It is common practice to use intensity standards (usually introducing the same objects of known attenuation into all scans) when carrying out XMT measurements, and this is usually sufficient for samples containing large features with high contrast. In these cases, variations in machine behaviour or beam energy over time (which is equivalent to variations in the threshold value) will be small. However for small objects, especially in low contrast materials or when volume changes can alter the bulk attenuation along the beam path, simple intensity calibration is unlikely to be sufficient. In this case we recommend having both volume and intensity references, especially for smaller grains. The number of reference features needs to be sufficient for suitably accurate volume determination, and the features should not change between scans over a time series experiment. In our particular example of grain dissolution, an appropriate standard could consist of an unaltered particle of the ore that is present in all scans. Ideally this procedure will be carried

283 out using a phantom containing a sufficient number of features for which the individual volumes  
284 are known, as this will allow not only consistent, but also accurate results. In this specific example  
285 the volumes in the reference image are not known and thus it is only consistency that is achieved  
286 by using this method.

287 The reason why the correction of systematic errors has been left to last is that it is important to  
288 know which errors or discrepancies can be eliminated by appropriate adjustment of the thresholds  
289 and which errors are random.

290 The relative difference in volume ( $\Delta V/V$ ) of the grains in two independent scans of the same  
291 volume collected under identical machine settings and analysed using the same thresholding  
292 algorithm (maximum entropy) should be negligible and yet plotting the  $\Delta V/V$  as a function of grain  
293 volume shows a systematic error in the volume, especially for smaller grains (Figure 8a). The  
294 discrepancy between the 2 scans will contain both systematic and random components and  
295 therefore the random component is reduced by combining measurements from 100 similar sized  
296 grains. This virtually eliminates the random component of the error for the larger particles, but it is  
297 still significant for the smaller particles (below about 100 voxels).

298 It is expected that much of the systematic difference between the images will be caused by a small  
299 inconsistency in the threshold value and it is the smallest grains that have the largest discrepancy.  
300 The same equation form that fitted the smaller particles in Figure 4 is therefore fitted to this data,  
301 namely a power law relationship with an exponent of  $-2/3$ .

302 Since the expected standard deviation in the average of the 100 grains used to generate each of  
303 the points in Figure 8 is known from Figure 6, the 95% confidence interval for the fitted equation  
304 can be plotted. For the smaller grains virtually all the points fit within this confidence interval,  
305 which would be expected if the assumed form for the data is correct. The difference in volume for  
306 the larger grains lies outside the confidence interval, but the power law relationship with an  
307 exponent of  $-2/3$  is only expected to fit the data for the smaller particles.

308 If there were no systematic error there should be no trend in the discrepancy and the data should  
309 be scattered around zero, with a larger scatter at smaller sizes. To try and achieve this, the  
310 correction to the threshold required to eliminate the systematic error can be estimated based on  
311 the prefactor  $k$  in the fitted power law (Figure 8a). The required change in threshold that  
312 corresponds to this value of  $k$  can be obtained from Figure 5. In this case an increase in the relative  
313 threshold value of about 1.5% was required<sup>2</sup>. The power law relationship between the change in  
314 volume and the volume when the threshold is adjusted means that even this small change has  
315 quite a large effect on the smallest grains. If this change in threshold has the same relative effect

---

<sup>2</sup> The prefactor in the power law fit to the data in Figure 8 (0.3344) is used to read off the required shift in threshold from Figure 5 (1.5%). Note that a positive prefactor implies that the measured volume is too large and that the threshold for the mineral grains must thus be increased.

on the measured volume in the subsequent image as it did on the reference image, then the systematic error should be eliminated.

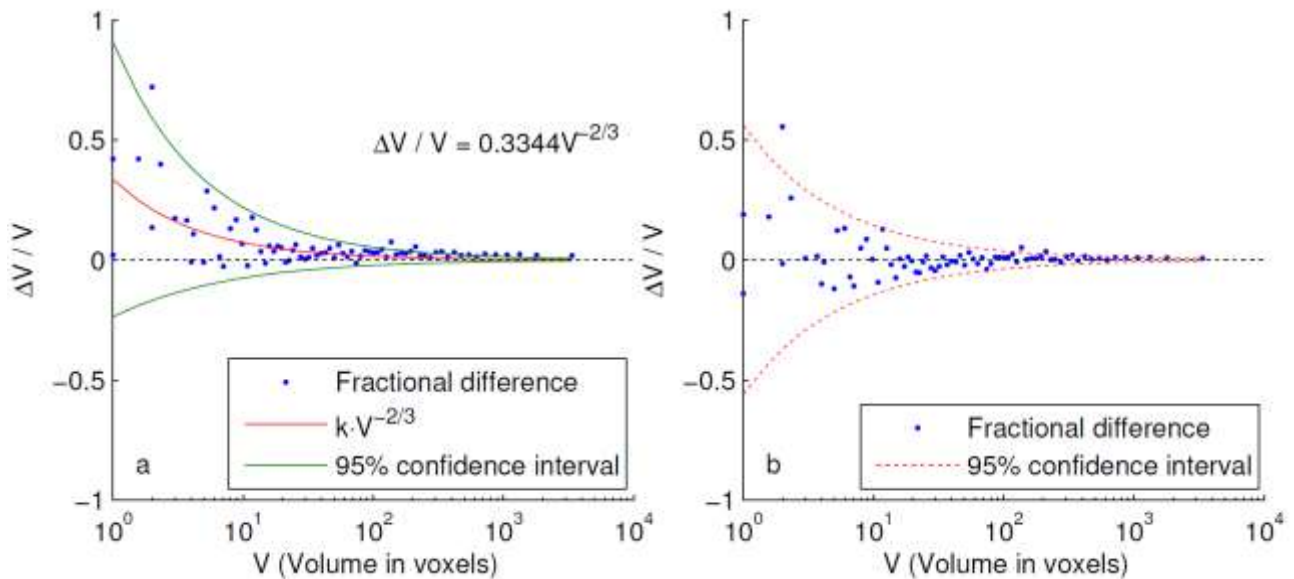


Figure 8. The discrepancy in grain volumes between: a) a reference scan and a repeated scan of the same ore particle before. b) after threshold correction (+1.5% in relative threshold value).

Figure 8b shows the discrepancy in the volume once this change in threshold has been applied. It can be seen from the 95% confidence intervals that this correction has resulted in discrepancies in volume that are consistent with no systematic error in the size of the smallest grains. While the correction was not based upon the size of the largest grains, the systematic error in their size was reduced from about 2% before correction to 0.8% after correction.

The correction has virtually no impact on the random component of the error, with the relationship between the standard deviation in the measured volume and volume itself for the corrected and uncorrected data being virtually the same (see Figure 6). This indicates that the random and systematic errors are independent of one another in this system. It also means that the random error can be accurately assessed without using a size standard as this component of the error is very insensitive to the specific threshold value used.

## 5. Conclusions

This paper described methodologies for quantifying, and correcting for, both the systematic and random contributions to the uncertainty and error in the measurement of the volume of objects using XMT. In particular, it showed the strong dependency on volume relative to voxel size that these errors have. To achieve a desired level of uncertainty due to random errors, the results from repeat scans or scans of similar objects need to be combined. The paper showed how many objects of a given size needed to be combined, providing a guide for future studies. For instance a single object of a thousand voxels has an uncertainty of 5% in its volume, while 12 objects of a hundred voxels would need to be combined to achieve the same level of uncertainty. A

341 methodology for eliminating systematic errors based on knowledge of how changes in threshold  
342 effect the measured volume and its dependency upon size was also developed.

### 343 **Acknowledgements**

344 This study was performed in the Rio Tinto Centre for Advanced Mineral Recovery at Imperial  
345 College London, and at the Manchester X-ray Imaging Facility, which was funded in part by the  
346 EPSRC (EP/I02249X/1). The authors gratefully acknowledge Rio Tinto for their financial support for  
347 this project.

### 348 **References**

- 349 Aydoğan, N.A., Ergün, L., Benzer, H., 2006. High pressure grinding rolls (HPGR) applications in the  
350 cement industry. *Minerals Engineering* 19, 130-139.
- 351 Ghorbani, Y., Becker, M., Mainza, A., Franzidis, J.-P., Petersen, J., 2011. Large particle effects in  
352 chemical/biochemical heap leach processes – A review. *Minerals Engineering* 24, 1172-1184.
- 353 Gonzalez, R.C., Woods, R.E., Eddins, S.L., 2003. *Digital Image Processing Using MATLAB*. Prentice-  
354 Hall, Inc.
- 355 Kapur, J.N., Sahoo, P.K., Wong, A.K.C., 1985. A new method for gray-level picture thresholding  
356 using the entropy of the histogram. *Computer Vision, Graphics, and Image Processing* 29, 273-285.
- 357 Ketcham, R.A., 2006. Accurate Three-dimensional Measurements of Features in Geological  
358 Materials from X-ray Computed Tomography Data, *Advances in X-ray Tomography for*  
359 *Geomaterials*. ISTE, pp. 143-148.
- 360 Ketcham, R.A., Carlson, W.D., 2001. Acquisition, optimization and interpretation of X-ray  
361 computed tomographic imagery: applications to the geosciences. *Computers & Geosciences* 27,  
362 381-400.
- 363 Ketcham, R.A., Hildebrandt, J., 2014. Characterizing, measuring, and utilizing the resolution of CT  
364 imagery for improved quantification of fine-scale features. *Nuclear Instruments and Methods in*  
365 *Physics Research Section B: Beam Interactions with Materials and Atoms* 324, 80-87.
- 366 Otsu, N., 1979. A threshold selection method from gray-level histograms. *IEEE Transactions on*  
367 *Systems, Man and Cybernetics* 9, 62-66.
- 368 Puncreobutr, C., Lee, P.D., Hamilton, R.W., Phillion, A.B., 2012. Quantitative 3D Characterization of  
369 Solidification Structure and Defect Evolution in Al Alloys. *JOM* 64, 89-95.
- 370 Stock, S.R., 1999. X-ray microtomography of materials. *International Materials Reviews* 44, 141-  
371 164.
- 372 Studholme, C., Hill, D.L.G., Hawkes, D.J., 1999. An overlap invariant entropy measure of 3D medical  
373 image alignment. *Pattern Recognition* 32, 71-86.

374 Yue, S., Lee, P.D., Poologasundarampillai, G., Jones, J.R., 2011. Evaluation of 3-D bioactive glass  
375 scaffolds dissolution in a perfusion flow system with X-ray microtomography. *Acta Biomaterialia* 7,  
376 2637-2643.

377

378